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Synthesis, Phase Evolutions, and Properties of the Perovskite Ba₆Ca₂Ti₂Nb₃O_{19.5}

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Abstract

In the present study, $Ba_6Ca_2Ti_2Nb_{2+x}O_{17+2.5x}$ (BCTN, $x \le 1.0$) materials were synthesized with the traditional solid-state reaction method. The phase evolution, structure, relative permittivity, and electrical conduction properties of these materials were studied. The results revealed that increasing the Nb content and improving the calcining temperature favored the formation of a pure simple cubic perovskite structure. For the end member of $Ba_6Ca_2Ti_2Nb_{2+x}O_{17+2.5x}$, i.e. $Ba_6Ca_2Ti_2Nb_3O_{19.5}$, Rietveld structural refinement disclosed that about 25 % Ca ions and all the Ba ions occupied the A site. While the remaining cations were disordered and occupied the B site. Thus, it can also be written as $(Ba_6Ca_{0.5})(Ca_{1.5}Ti_2Nb_3)O_{19.5}$. This material exhibited high permittivity up to 80 and the electrical conductivity was 1.33×10^{-4} S/cm at 800 °C which was dominated by electronic conduction.

Keywords: $Ba_6Ca_2Ti_2Nb_3O_{19.5}$, simple cubic perovskite, B-site disorder, high permittivity, alternating current impedance.

I. Introduction

Perovskite oxides, due to their diverse structural features and interesting physical and chemical properties, are important in current solid-state chemistry and physics, and have wide application 1,2. Thanks to their high relative permittivity, which is desirable for minimization of the microwave circuit component ^{3,4}, ABO₃-type perovskite materials have been very attractive for the electronics industry over the past 30 years. The cubic ABO3 perovskite structure consists of a three-dimensional array of cornersharing BO₆ octahedra, with the larger A cations in a 12coordinate site, and arises from ordered B cation occupancy of 25 % of the octahedral sites between close-packed AO₃ layers stacked along the <111> direction of the cubic unit cell, as shown in Fig. 1. Cation site order in perovskite-derived structures is important, as it can control the physical properties such as dielectric loss 5,6, magnetoresistance ⁷, magnetic order ^{8,9}, and ionic mobility ^{10,11}. The cation order is of particular importance for controlling microwave dielectric properties 5,12, such as the permittivity at microwave frequencies.

In 2008, Kuang et al. ¹ reported a new six-layer perovskite-related structure Ba₆Na₂Nb₂M₂O₁₇ (M = P, V), which consists of cubic (c) BaO₃ layers and oxygen-deficient pseudocubic (c') BaO₂ layers stacked in the sequence c'cccc. In Ba₆Na₂Nb₂M₂O₁₇, two-dimensional slabs of the well-known 2:1 octahedral cation-ordered perovskite motif are isolated between layers of tetrahedral units formed by anion vacancy ordering: two consecutive NbO₆ octahedral layers are sandwiched by two single NaO₆ octahedral layers, which, in turn, connect with two isolated MO₄ tetrahedral layers. Both oxides are derived from the 2:1 ordered perovskite structure (e.g. $Ba_3ZnTa_2O_9$) ¹³ by the ordered removal of O atoms in every sixth BaO_3 layer. Both materials exhibit a relative permittivity of ~ 20-23.

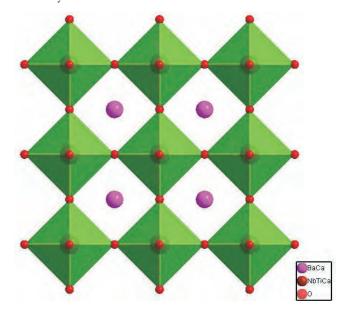


Fig. 1: The crystal structure of a typical simple cubic perovskite.

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In this work, we originally intended to replace the Na and P/V cations with Ca and Ti, respectively, to prepare the Ba₆Ca₂Ti₂Nb₂O₁₇ material with the hexagonal perovskite structure that is iso-structural with Ba₆Na₂Nb₂M₂O₁₇. However, the secondary cubic perovskite phase in the products could not be avoided, and increasing the Nb content would ultimately yield pure simple cubic perovskite Ba₆Ca₂Ti₂Nb₃O_{19.5} as a product. This novel simple cubic perovskite material ceramicshowed high permittivity ($\varepsilon_r \sim 80$) and impedance data revealed the dominating electronic conduction, without ionic conduction being observed.

II. Experimental procedure

Ba₆Ca₂Ti₂Nb_{2+x}O_{17+2.5x} (BCTN) materials were prepared based on a routine solid-state reaction, using highpurity BaCO₃ (99.99 %, Damao Chemical Reagent Factory), CaCO₃ (99.99 %; Aladdin), TiO₂ (99.99 %; Aladdin), Nb₂O₅ (99.99 %, Aladdin) as starting materials, which were dried at 400 °C for 3 hours before they were weighed. The starting materials were weighed stoichiometrically, mixed in ethanol with an agate mortar, dried at room temperature and precalcined at 1 200 °C for 12 hours in an alumina crucible with heating and cooling rates of 5 K/min. The pre-calcination process was employed to decompose the carbonate CaCO₃ and BaCO₃ into CaO, BaO, CO₂ and improve homogeneity of the mixture, which would be beneficial to the formation of pure phase product in the final sintering. The precalcined powder was then mixed with 5 % polyvinyl alcohol (PVA) solution as an organic binder and pressed into pellets under 140 MPa pressure. The pellets were buried in sacrificial powders of the same composition in aluminum crucibles covered with lids and heated at 600 °C for 100 min to remove the PVA, and then fired at 1 400 °C - 1 600 °C for 12 - 24 hours with heating rates of 5 K/min and cooling rates 0.5 K/min. The densities of the pellets were calculated by using the geometric size (diameter and thickness) and the mass of the pellets.

The phase compositions were checked by means of the powder X-ray diffraction (XRD) technique, which was performed on a D8 ADVANCE powder diffractometer with CuK α radiation. The XRD data were collected over a 2θ range of $10-80^\circ$. The powder XRD data was analyzed with Rietveld refinement using Topas Academic software. The microstructure of the pellets was examined using a Quanata 400F scanning electron microscope (SEM). Prior to the SEM experiment, gold was sprayed onto the surface to form a thin conducting layer.

AC impedance spectroscopy measurement was carried out on a Solartron 1260A impedance/gain-phase analyzer with the frequency range from 10⁷ Hz to 10⁻¹ Hz over a temperature range from room temperature to 800 °C. For these measurements, the pellets were coated with gold paste and fired at 550 °C for 90 min to remove the organic components in order to form electrodes. In order to obtain accurate permittivity of the material itself, the blank contribution arising from the sample holder and connection cables, measured on an open circuit without a pellet, was subtracted from the measured capacitance of the pellets. The microwave dielectric properties were measured

with the Hakki-Coleman dielectric resonator method ³ with the TE₀₁₁ mode using an Agilent N5230A network analyzer.

III. Results and Discussion

(1) Phase evolution of $Ba_6Ca_2Ti_2Nb_{2+x}O_{17+2.5x}$

As reported in Kuang *et.al.*'s previous work, hexagonal perovskite Ba₆Na₂Nb₂M₂O₁₇ (M = P, V) could be obtained with a traditional solid-state reaction method. Herein, we originally intended to completely replace the Na and M in Ba₆Na₂Nb₂M₂O₁₇ with Ca and Ti, respectively, to prepare the hexagonal perovskite Ba₆Ca₂Ti₂Nb₂O₁₇. However, mixed hexagonal and cubic perovskite phases were always there in the products, no single phase could be obtained, and changing the preparation conditions, such as increasing the temperature or prolonging the dwell time, did not work. An exemplary XRD pattern for the product of Ba₆Ca₂Ti₂Nb₂O₁₇ prepared at 1500 °C for 12 hours can be seen in Fig. 2.

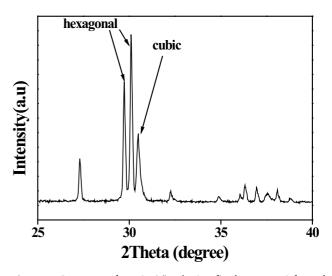


Fig. 2: XRD pattern of $Ba_6Ca_2Ti_2Nb_2O_{17}$ fired at 1 500 °C for 12 h.

Interestingly, based on the analysis of quantities of experiments, we found that increasing the Nb content would unexpectedly increase the cubic phase proportion in the products, and even a single cubic phase product could be yielded for the composition Ba₆Ca₂Ti₂Nb₃O_{19.5}, as shown in Fig. 3a. To investigate the exact cation distribution in the unit cell of Ba₆Ca₂Ti₂Nb₃O_{19.5}, Rietveld structural refinements¹⁴ against the XRD data were carried out using the simple cubic perovskite-type BaTiO₃ (space group of Pm3m) as a structural model, with one A site (1a for Ba), one B site (1b for Ti), and one C site (3c for oxygen). The refinement confirmed that besides all the Ba ions, 1/4 of the Ca ions in Ba₆Ca₂Ti₂Nb₃O_{19.5} are also located in the A site, while the remaining Ca ions and all the Ti and Nb ions are randomly distributed in the B site. The Rietveld fitting plots of the XRD data for Ba₆Ca₂Ti₂Nb₃O_{19.5} are shown in Fig. 3b. The final refined structural parameters of Ba₆Ca₂Ti₂Nb₃O_{19.5} are listed in Table 1.

Table 1: The final refined structural parameters for Ba₆Ca₂Ti₂Nb₃O_{19.5}.

Atoms	Site	\boldsymbol{x}	\mathcal{Y}	z	Occupancy	$B_{iso}(\mathring{A}^2)$
Ba	1a	0	0	0	0.923(1)	1.61(4)
Ca1	1a	0	0	0	0.077(1)	1.61(4)
Ca ₂	1b	0.5	0.5	0.5	0.231(1)	1.15(3)
Ti	1b	0.5	0.5	0.5	0.307(2)	1.15(3)
Nb	1b	0.5	05	0.5	0.462(1)	1.15(3)
O	3c	0	0.5	0.5	1	2.36(2)

Yobs (b) (a) Yeale **Bragg positions** Intensity(a.u.) Ydiff Intensity(a.u) PDF#74-1961 BaTiO₃ 20 30 40 50 70 10 60 30 40 50 10 20 60 70 80

Fig. 3: (a) XRD data and (b) corresponding Rietveld fitting plots of Ba₆Ca₂Ti₂Nb₃O_{19.5}.

2Theta(degree)

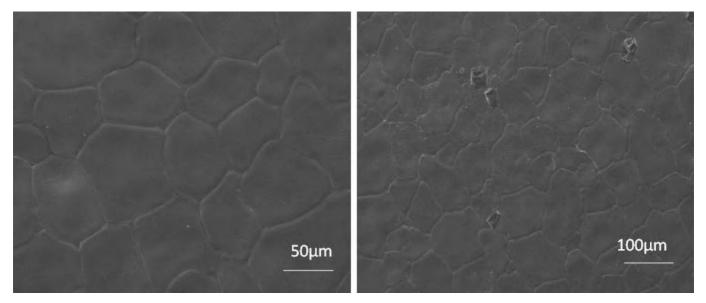


Fig. 4: SEM images of surface morphologies of the Ba₆Ca₂Ti₂Nb₃O_{19.5} pellets.

The microstructure of the $Ba_6Ca_2Ti_2Nb_3O_{19.5}$ pellet fired at 1 500 °C for 12 hours was then examined by means of SEM, as shown in Fig. 4. The images showed a dense microstructure free of pores, which is consistent with its relative density of ~93 %, which was calculated based on the geometric size (diameter and thickness) and the mass

of the pellet. The grains sizes are mainly in the range of $50-100 \,\mu m$ with well-defined grain boundaries.

2Theta(degree)

AC impedance spectroscopies ¹⁵ were then recorded on the dense pellet of the Ba₆Ca₂Ti₂Nb₃O_{19.5} sample. The impedance showed two separated arcs within the whole measured temperature range, as can be seen from Fig. 5

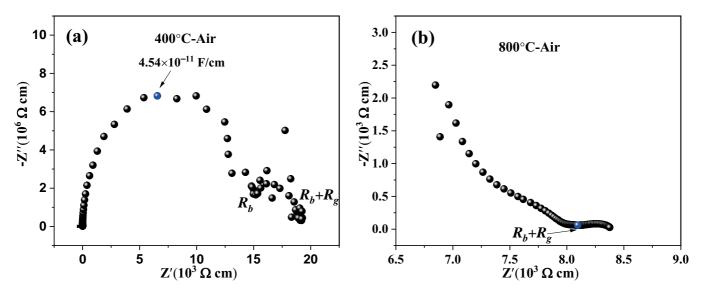
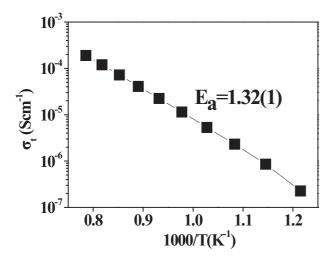


Fig. 5: Impedance plots of Ba₆Ca₂Ti₂Nb₃O_{19.5} measured at 400 and 800 °C.



 $\textbf{Fig. 6:} Arrhenius plots of bulk conductivity for Ba_6 Ca_2 Ti_2 Nb_3 O_{19.5}.$

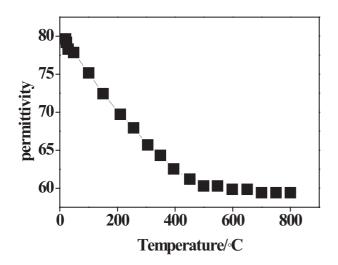


Fig. 7: Permittivity of ${\rm Ba_6Ca_2Ti_2Nb_3O_{19.5}}$ pellets collected from 25 °C – 800 °C.

for the impedance plots recorded at 400 °C and 800 °C. The arcs in the high- and low-frequency range correspond-

ed to the grain and grain boundary response, respectively. No electrode response was observed, which excluded the ionic conduction in this material. The total resistivity of the $\rm Ba_6Ca_2Ti_2Nb_3O_{19.5}$ pellet was simply estimated as the intercept of the semicircular arc at low frequency. The Arrhenius plot of total conductivities for the $\rm Ba_6Ca_2Ti_2Nb_3O_{19.5}$ pellet is shown in Fig. 6. The conductivity of $\rm Ba_6Ca_2Ti_2Nb_3O_{19.5}$ ceramic varied within 2.53×10^{-7} - 1.33×10^{-4} S/cm in 300-800 °C.

The dielectric constants (ϵ_r) obtained from the impedance spectroscopies showed a decrease from ϵ_r ~80 at room temperature to ϵ_r ~60 at 800 °C (Fig. 7). The temperature coefficient of permittivity ranged in $10^{-3} - 10^{-5}$ from room temperature to 400 °C, as shown in Fig. 8.

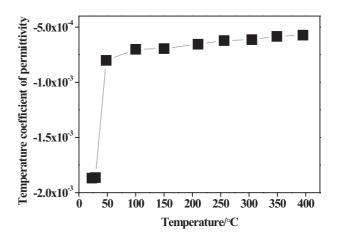


Fig. 8: The temperature coefficient of permittivity of Ba₆Ca₂Ti₂Nb₃O_{19.5}.

IV. Conclusions

The Ba₆Ca₂Ti₂Nb_{2+z}O_{17+ δ} (BCTN) materials were synthesized with a traditional solid-state reaction method. The phases, structures, and electrical properties were studied. The results revealed that the Nb content increase fa-

vored the formation of pure simple cubic-perovskite-structured $Ba_6Ca_2Ti_2Nb_3O_{19.5}$. The Rietveld structural refinement showed that about 25 % of Ca ions and all the Ba ions occupied the A site, while the remaining cations were disordered and occupied the B site. The electrical conductivity of this material was 2.52×10^{-7} - 1.23×10^{-4} S/cm within the temperature range of 300-800 °C and dominated by electronic conduction. The dielectric constants decreased from 80 at room temperature to ~60 at 800 °C.

Acknowledgments

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